

Absorbed Dose in Radioactive Media

Chapter 5

F.A. Attix, Introduction to Radiological Physics and Radiation Dosimetry

Outline

- General dose calculation considerations, absorbed fraction
- Radioactive disintegration processes and associated dose deposition
 - Alpha disintegration
 - Beta disintegration
 - Electron-capture transitions
 - Internal conversion
- Summary

Introduction

- We are interested in calculating the absorbed dose in radioactive media, applicable to cases of
 - Dose within a radioactive organ
 - Dose in one organ due to radioactive source in another organ
- If conditions of CPE or RE are satisfied, dose calculation is straightforward
- Intermediate situation is more difficult but can be handled at least in approximations

Radiation equilibrium

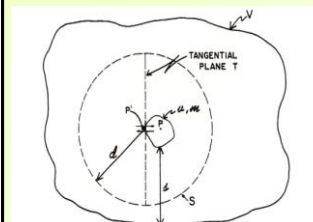


FIGURE 4-1. Radiation equilibrium. Extended volume F contains a homogeneous medium and a homogeneous isotropic source distribution. Radiation equilibrium will exist in the smaller internal volume s if the maximum distance of penetration (d) of primary rays plus their secondaries is less than the minimum separation (s) of s from the boundary of V . Neutrinos are ignored. (See text.)

- a. The atomic composition of the medium is homogeneous
- b. The density of the medium is homogeneous
- c. The radioactive source is uniformly distributed
- d. No external electric or magnetic fields are present

Charged-particle equilibrium

- Each charged particle of a given type and energy leaving the volume is replaced by an identical particle of the same energy entering the volume
- Existence of RE is sufficient condition for CPE
- Even if RE does not exist CPE may still exist (for a very large or a very small volume)

Limiting cases

- Emitted radiation typically includes both photons (longer range) and charged particles (shorter range)
- Assume the conditions for RE are satisfied
 - Consider two limited cases based on the size of the radioactive object

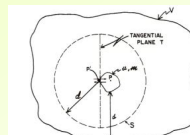


FIGURE 4-1. Radiation equilibrium. Extended volume F contains a homogeneous medium and a homogeneous isotropic source distribution. Radiation equilibrium will exist in the smaller internal volume s if the maximum distance of penetration (d) of primary rays plus their secondaries is less than the minimum separation (s) of s from the boundary of V . Neutrinos are ignored. (See text.)

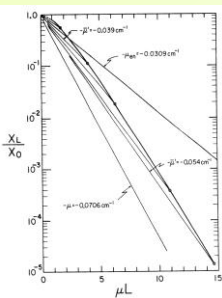
Limiting cases: small object

- A radioactive object V having a mean radius not much greater than the maximum charged-particle range d
- CPE is well approximated at any internal point P that is at least a distance d from the boundary of V
- If $d \ll 1/\mu$ for the γ -rays, the absorbed dose D at P approximately equals to the energy per unit mass of medium that is given to the charged particles in radioactive decay (less their radiative losses)
- The photons escape from the object and are assumed not to be scattered back by its surroundings

Limiting cases: large object

- A radioactive object with mean radius $\gg 1/\mu$ for the most penetrating γ -rays
- RE is well approximated at any internal point P that is far enough from the boundary of V so γ -ray penetration through that distance is negligible
- The dose at P will then equal the sum of the energy per unit mass of medium that is given to charged particles plus γ -rays in radioactive decay

Limiting cases: large object



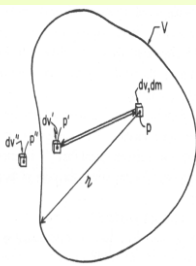
- Deciding upon a maximum γ -ray "range" for this case requires quantitative criterion
 - For primary beam (μ) only less than 0.1% through 7 mean free paths
 - Taking into account scatter, (broad beam geometry, $\bar{\mu}'$) will increase the required object size to satisfy the attenuation objectives (e.g. 16 mean free paths to achieve 0.1% penetration)

Absorbed fraction

- An *intermediate-size* radioactive object V
- Dose at P will then equal the sum of the energy per unit mass of medium that is given to charged particles plus dose from γ -rays
- To estimate the dose from γ -rays define absorbed fraction:

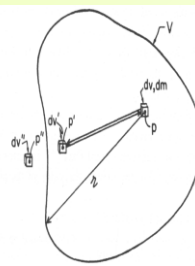
$$AF = \frac{\gamma\text{-ray radiant energy absorbed in target volume}}{\gamma\text{-ray radiant energy emitted by source}}$$

Intermediate-size radioactive object



- Consider volume V filled by a homogeneous medium and a uniformly distributed γ -source
- The volume may be surrounded by
 - Case 1: infinite homogeneous medium identical to V , but non-radioactive (an organ in the body)
 - Case 2: Infinite vacuum (object in air)

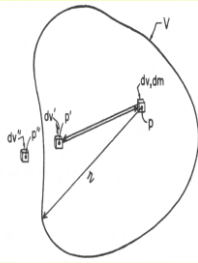
Case 1



- Reciprocity theorem: energy spent in dv due to the source in dv' :

$$\mathcal{E}_{dv',dv} = \mathcal{E}_{dv,dv'} \text{ and } \mathcal{E}_{dv,V} = \mathcal{E}_{V,dv}$$
- No source in dv''
- Define \mathcal{R}_{dv} as the expectation value of the γ -ray radiant energy emitted by the source in dv , and $\bar{\epsilon}_{dv,V}$ the part of that energy that is spent in V , then the absorbed fraction with respect to source dv and target V is

Case 1 continued



- The absorbed fraction with respect to source dv and target V is :

$$AF_{dv,V} = \frac{\bar{\mathcal{E}}_{dv,V}}{R_{dv}} = \frac{\bar{\mathcal{E}}_{V,dv}}{R_{dv}}$$

source, target

- Estimates reduction in absorbed dose relative to RE condition
- For very small radioactive objects ($V \rightarrow dv$) this absorbed fraction approaches zero; for an infinite radioactive medium it equals unity

Case 1 continued

- Reduction in the absorbed dose due to γ -rays energy escaping from V
- Can estimate AF using mean effective attenuation coefficient $\bar{\mu}'$ for γ -rays energy fluence through a distance r in the medium

$$AF_{dv,V} = \frac{1}{4\pi} \int_{\theta=0}^{\pi} \int_{\beta=0}^{2\pi} (1 - e^{-\bar{\mu}'r}) \sin \theta d\theta d\beta$$

- For poly-energetic sources have to find an average value of the absorbed fraction

Case 1 example

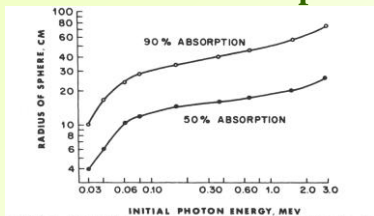
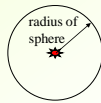


FIGURE 5.2. Radius of unit-density tissue sphere needed to absorb 50% and 90% of the emitted photon energy from a central point source in an infinite homogeneous medium. (After Brownell et al., 1968.) Reproduced with permission of the authors and The Society of Nuclear Medicine.

- Dose calculations published in MIRD reports
- The larger the radius, the lower the energy – the closer to RE condition ($AF=1$)



Case 2

Radioactive object V surrounded by vacuum

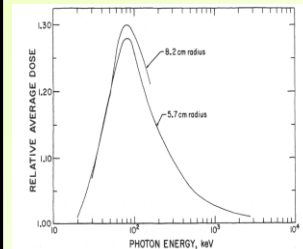


FIGURE 5.3. Ratio of average absorbed doses in uniformly radioactive tissue spheres with/without surrounding non-radioactive tissue medium. Lower curve: 780-g (5.7-cm-radius) sphere; upper curve: 2.5-kg (8.2-cm-radius) sphere. (Data from Elliott, 1968.)

- More difficult to calculate
- The reciprocity theorem is only approximate due to the lack of backscattering
- Dose is lower than in Case 1

Case 2

- To obtain a crude estimate of the dose at some point P within a uniformly γ -active homogeneous object, it may suffice to obtain the average distance \bar{r} from the point to the surface of the object by

$$\bar{r} = \frac{1}{4\pi} \int_{\theta=0}^{\pi} \int_{\beta=0}^{2\pi} r \sin \theta d\theta d\beta$$

- Then one may employ $\mu_{en} = \bar{\mu}'$ in the straight-ahead approximation to obtain

$$AF_{dv,V} \cong 1 - e^{-\mu_{en}\bar{r}}$$

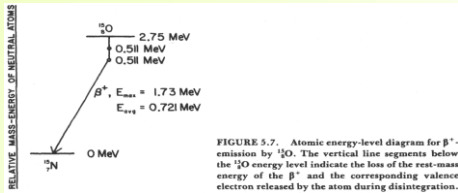
Radioactive disintegration processes

- Radioactive nuclei tend to undergo transformations to a more stable state through expulsion of energetic particles

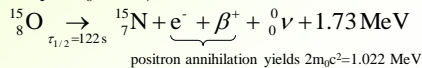
	ΔZ	ΔA
α -particle	-2	-4
β^- -particle†	+1	0
β^+ -particle†	-1	0
γ -ray	0	0

- The total mass, energy, momentum, and electric charge are conserved
- Energy equivalent of the rest mass:
 - 1 amu=1/12 of the mass of $^{12}_6\text{C}$ nucleus=931.50 MeV
 - 1 electron mass=0.51100 MeV

Beta disintegration



- In β^+ disintegration valence electron is emitted
- Example: $^{15}_8\text{O} \rightarrow ^{15}_7\text{N}$



Absorbed dose from β -disintegration

- Under CPE condition $D=nE_{\text{avg}}$ MeV/g for n disintegration per gram of medium
- Any additional contributions to energy deposition due to γ -rays must be included for RE condition
- Radiative losses by β -rays, such as bremsstrahlung and in-flight annihilation, are ignored

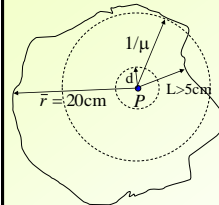
Example 5.1

- Uniformly distributed β - and γ -ray source
- The rest-mass loss is spent
 - half in 1 -MeV γ -ray production and
 - half in β^- -decay, for which $E_{\text{max}}=5 \text{ MeV}$ and $E_{\text{avg}}=2 \text{ MeV}$
- The point of interest P is located $>5 \text{ cm}$ inside the boundary of the object, at an average distance $\bar{r} = 20 \text{ cm}$ from the boundary.
- $\mu_{\text{en}}=0.0306 \text{ cm}^{-1}$ and $\mu = 0.0699 \text{ cm}^{-1}$ for the γ -rays
- A total energy of 10^{-2} J converted from rest mass in each kg of the object
- Estimate the absorbed dose at P

Example 5.1

- For β -ray $E_{\text{max}}=5 \text{ MeV}$ corresponds to maximum particle range (Appendix E) $d \sim 2.6 \text{ cm} \ll 1/\mu=14.3 \text{ cm}$
- CPE exists at P , therefore dose due β -rays:

$$D_{\beta} = E_{\text{avg}} \times \frac{1}{2} \times 10^{-2} \text{ J/kg} = 10^{-2} \text{ J/kg}$$



- For γ -ray 20 cm is not $\gg 1/\mu=14.3 \text{ cm}$
- RE does not exist at P , therefore have to use absorbed fraction:

$$AF \cong 1 - e^{-\mu_{\text{en}} \bar{r}} = 0.46$$

$$D_{\gamma} = 0.46 \times (1/2) \times 10^{-2} \text{ J/kg} = 2.3 \times 10^{-3} \text{ J/kg}$$

$$D_{\text{tot}} = (2.3 + 10) \times 10^{-3} \text{ J/kg} = 1.23 \times 10^{-2} \text{ J/kg}$$

Example 5.2

- What is the absorbed dose rate (Gy/h) at the center of a sphere of water 1 cm in radius, homogeneously radioactivated by $^{32}_{15}\text{P}$, with 6×10^5 disintegrations per second occurring per gram of water? (Assume time constancy.)

Example 5.2

- $E_{\text{max}}=1.71 \text{ MeV}$ corresponds to maximum particle range of $\sim 0.8 \text{ cm} < 1 \text{ cm}$
- CPE condition
- Absorbed dose rate: $\dot{D} = \dot{N} \times E_{\text{avg}}$

$$\begin{aligned} \dot{D} &= 6 \times 10^5 \frac{\text{dis}}{\text{g sec}} \times 0.694 \frac{\text{MeV}}{\text{dis}} \\ &= 4.164 \times 10^5 \frac{\text{MeV}}{\text{g sec}} \times 3600 \frac{\text{s}}{\text{hr}} \times 1.602 \times 10^{-10} \frac{\text{Gy}}{\text{MeV/g}} \\ &= 2.4 \text{ Gy/h} \end{aligned}$$

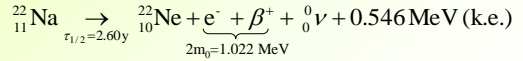
Electron-capture transitions

- Parent nucleus captures its own atomic electron from K-shell (~90% probability) or L-shell (~10% probability) and emits monoenergetic neutrino
- Resulting shell vacancy is filled with electron from a higher orbit, leading to emission of a fluorescence x-ray
- Process competing with β^+ disintegrations

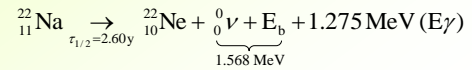
Electron-capture transitions

- Example: $^{22}_{11}\text{Na} \rightarrow ^{22}_{10}\text{Ne}$ with half-life for both branches of 2.60 years

– β^+ branch



– EC branch



Electron-capture transitions

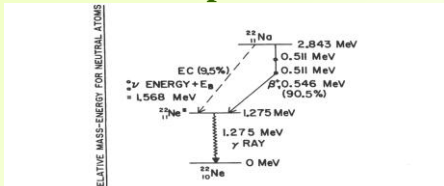


FIGURE 5.8. Atomic energy-level diagram for β^+ and EC disintegration of $^{22}\text{Na} \rightarrow ^{22}\text{Ne}$. ($E_{\beta_{\text{max}}} = 0.546\text{ MeV}$; $E_{\beta^+} = 0.216\text{ MeV}$.)

- Binding energy for K-shell $E_b \sim 1\text{ keV}$
- For electron capture to occur the minimum atomic mass decrease of $2m_0$ between the parent and daughter nuclei is required to supply β^+ with kinetic energy

Absorbed dose for EC process

- Most of the energy is carried away by neutrino
- The only available energy for dose deposition comes from electron binding term E_b , which is very small compare to that of neutrino

Internal conversion

- An excited nucleus can impart its energy directly to its own atomic electron, which then escapes with the net kinetic energy of $h\nu - E_b$ ($h\nu$ is the excitation energy)
- No photon is emitted in this case
- Process competing with γ -ray emission
- Internal conversion coefficient is the ratio of N_e/N_γ

Internal conversion

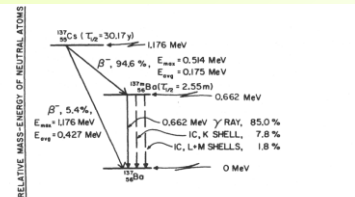


FIGURE 5.9. Atomic energy-level diagram for $^{137}\text{Cs} \rightarrow ^{137}\text{Ba}$, illustrating competition between γ -ray emission and internal conversion. Percentages all refer to disintegrations of parent atoms of ^{137}Cs . ($\alpha_{\text{K}} = 0.0916$, $K/(L + M + \dots) = 4.41$.)

- Example: $^{137}_{55}\text{Cs} \rightarrow ^{137}_{56}\text{Ba}$

Absorbed dose for internal conversion

- If IC occurs in competition with γ -ray emission, it results in increase in absorbed dose in small objects (CPE condition) due to release of electron locally depositing the energy

$$E_{IC} = hv - E_b$$

- In addition electron binding energy is contributed to the dose unless it escapes as a fluorescence x-ray

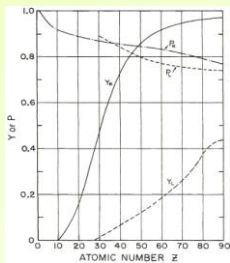
Absorbed dose for internal conversion

- If the fraction $p = 1 - AF$ of these fluorescence x-rays escape, then the energy contributed to dose per IC event under CPE condition

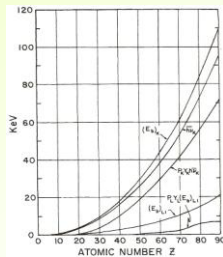
$$f_{IC} = hv - p_K Y_K h\bar{v}_K - p_L Y_L h\bar{v}_L$$

- Using straight-ahead approximation $p \cong e^{-\mu_{en} r}$
- Values of fluorescence yield $Y_{K,L}$ and the mean emitted x-ray energies $h\bar{v}_{K,L}$ are tabulated

Fluorescence data



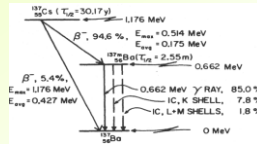
Fluorescence yield for K and L shells



Electron binding energies and mean fluorescence x-ray energies, K and L shells

Example 5.5

- A sphere of water 10 cm in diameter contains a uniform source of ^{137}Cs undergoing 10^3 disintegration per g s. What is the absorbed dose at the center, in grays, for a 10-day period, due only to the decay of $^{137\text{m}}_{56}\text{Ba}$? Use the mean-radius straight-ahead approximation.



For γ -ray of 0.662 MeV in water $\mu_{en}=0.0327 \text{ cm}^{-1}$

Example 5.5

- First check RE condition: $1/\mu=30.6 \text{ cm} \gg r=5$
- Find absorbed fraction $AF \cong 1 - e^{-0.0327 \times 5} = 0.151$
- Dose in 10 days = $8.64 \times 10^5 \text{ s}$

$$D_\gamma = 10^3 \frac{\text{dis}}{\text{g s}} \times 0.85 \frac{\gamma\text{-rays}}{\text{dis}} \times 0.662 \frac{\text{MeV}}{\gamma\text{-ray}} \times 1.602 \times 10^{-10} \frac{\text{Gy}}{\text{MeV/g}} \times 8.64 \times 10^5 \text{ s} \times AF = 1.17 \times 10^{-2} \text{ Gy}$$

Example 5.5

- For the K-shell conversion process we need $Y_K=0.90$, $h\bar{v}_K = 0.032 \text{ MeV}$, and $\mu_{en}=0.13 \text{ cm}^{-1}$ for 0.032 MeV. Then $p_K \cong e^{-\mu_{en} r} = e^{-0.13 \times 5} = 0.52$ and the dose contribution is

$$D_K = 10^3 \frac{\text{dis}}{\text{g s}} \times 0.078 \frac{\text{IC}(K)}{\text{dis}} \times (hv - p_K Y_K h\bar{v}_K) \frac{\text{MeV}}{\text{IC}(K)} \times 1.602 \times 10^{-10} \frac{\text{Gy}}{\text{MeV/g}} \times 8.64 \times 10^5 \text{ s} = 1.080(0.662 - 0.52 \times 0.90 \times 0.032) \times 10^{-2} \text{ Gy} = 6.99 \times 10^{-3} \text{ Gy}$$

Example 5.5

- Similarly, for the L+M+...-shell conversion process we need $Y_K=0.90$, $h\bar{\nu}_L = E_0^L = 6 \text{ keV}$, and $\mu_{\text{en}}=24 \text{ cm}^{-1}$ for 6 keV. Then $p_L \cong e^{-24 \times 5} = 0$ and the corresponding dose contribution

$$D_{ic}^L = 10^3 \frac{\text{dis}}{\text{g s}} \times 0.078 \frac{\text{IC}(L)}{\text{dis}} \times (h\nu - p_L Y_L h\bar{\nu}_L) \frac{\text{MeV}}{\text{IC}(L)}$$

$$\times 1.602 \times 10^{-10} \frac{\text{Gy}}{\text{MeV/g}} \times 8.64 \times 10^5 \text{ s}$$

$$= 1.65 \times 10^{-3} \text{ Gy}$$

- The total absorbed dose is $D_{\text{int}} = D_\gamma + D_{ic}^K + D_{ic}^L = 2.03 \times 10^{-2} \text{ Gy}$

Summary

- General approach to dose calculation within and outside of distributed radioactive source
- Radioactive disintegration processes and calculation of absorbed dose
 - Alpha disintegration
 - Beta disintegration
 - Electron-capture transitions
 - Internal conversion